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# Spectroscopic ellipsometry (SE) studies on vacuum-evaporated ZnSe thin films

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#### Abstract

Iodine-doped ZnSe thin films were prepared onto well-cleaned glass substrates using vacuum evaporation technique under a vacuum of  $3.4 \times 10^{-3}$  Pa. The composition, structural, optical and electrical properties of the deposited films were analyzed using Rutherford Backscattered Spectrometry (RBS), X-ray diffraction (XRD), spectroscopic ellipsometry (SE) and I–V characteristics. In the RBS analysis, the composition of the deposited film was calculated as (ZnSe)I<sub>0.003</sub>. The structure of the deposited film was identified as cubic, oriented along the (111) direction. The structural parameters such as particle size, strain and dislocation density values were calculated as 28.28 nm,  $1.38 \times 10^{-3} \ln^{-2} m^{-4}$  and  $1.29 \times 10^{15} \ln/m^2$ , respectively. Spectroscopic Ellipsometric (SE) measurements were done on the (ZnSe)I<sub>0.003</sub> thin films. The spectral data showed three distinct critical-point structures, including weak structures at  $E_0 + \Delta_0$ , indicating that the sample has a high crystalline quality. The optical band gap value of the deposited films was calculated as 2.63 eV using optical transmittance measurements. From the I–V characteristics studies, the conduction mechanism was found to be SCLC.

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## 1. Introduction

Zinc selenide (ZnSe) is well known as a high refractive index material in multilayer film combinations and as an infrared antireflection coating for solar cells, due to its wide band gap. The growth of II–VI compound semiconductors has attracted considerable attention due to their novel physical properties and wide range of applications in optoelectronic devices. Among these, zinc selenide (ZnSe) material with a direct band gap of 2.7 eV at room temperature, has potential application in light emitting diodes [1] and photoluminescence [2]. In photo-electronic and other properties of the II–VI class of compounds, thin films are highly structure sensitive which in turn can severely influence the device performance. ZnSe film is also used as a buffer layer in thin film solar cells [3]. A broad variety of techniques have been used for the deposition of ZnSe films, such as RF magnetron sputtering [4], metalorganic chemical vapour deposition (MOCVD) [5], molecular beam epitaxy (MBE) [6] and vacuum

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evaporation [7]. ZnSe has either a Sphalerite structure with lattice parameter a=5.6688 Å or a Wurtzite structure with lattice parameters a=3.820 Å and c=6.626 Å [8]. In this paper we report on the composition, structural and optical analysis of vacuum evaporated (ZnSe)I<sub>0.003</sub> thin films.

## 2. Experimental

High purity Zn (99.999%), Selenium (99.999%) and Iodine (99.99%) were used to prepare the iodine-doped ZnSe compound. Iodine-doped ZnSe compound was prepared by agitating the evacuated  $(3.4 \times 10^{-3} \text{ Pa})$  and sealed quartz ampoule containing the above mentioned elements at 875 °C. The prepared compound was light yellow in colour. (ZnSe)I<sub>0.003</sub> powder was thermally evaporated from a molybdenum boat onto well-cleaned glass substrates under a vacuum of  $3.4 \times 10^{-3}$  Pa. During evaporation, the substrates were maintained at room temperature. The composition of the film deposited onto a glass substrate was studied using Rutherford Backscattered spectrometry. 2 MeV He<sup>+</sup> ion beam was used for the experiment. Details of the experimental set-up can be found in Ref. [9]. The RBS analysis was performed using a simulation code Rump [10]. Thicknesses of the deposited films were measured by a Multiple Beam Interference (Fizeau fringes) technique [11] and the thickness of the as-deposited film was calculated as 225 nm. The surface morphology of the film was studied by SEM [JEOL, JSM 6400]. Structural analyses of the films were made by X-ray diffractormeter (SCINTAG; USA, Cu K<sub>a</sub> radiation at  $\lambda = 0.541$  nm) in the 2 $\theta$  range 20° to 70°. The optical transmittance spectra were recorded using a UV-VIS-NIR spectrophotometer in the wavelength range from 200-2500 nm.

### 2.1. Optical properties

The extinction coefficient 'k' at various wavelengths has been calculated using the equation [12]:

$$k = \frac{2.303\log\left(\frac{1}{T}\right)}{4\pi t} \tag{1}$$

where *T* is the transmittance and *t* is the thickness of the deposited film. The value of the absorption coefficient ( $\alpha$ ) can be evaluated from the formula [12]:

$$\alpha = \frac{4\pi k}{\lambda} \tag{2}$$

The relation between  $\alpha$  and incident photon energy  $h\nu$  can be written as [13]:

$$\alpha h v = C_1 \left( h v - E_g^d \right)^{1/2} \tag{3}$$

$$\alpha h v = C_2 (h v - E_{\sigma}^i)^2 \tag{4}$$

for direct allowed and indirect allowed transitions, respectively, where  $C_1$  and  $C_2$  are two constants,  $E_g^d$ ,  $E_g^i$  are the direct and indirect band gaps, respectively.

# 3. Results and discussion

### 3.1. Composition and surface analysis

The composition of the deposited films was measured using RBS analysis. The RBS spectrum for the as deposited film is shown in Fig. 1. The composition of the deposited films was calculated as  $(ZnSe)I_{0.003}$ /glass and it was found to be nearly stoichiometric with a negligible iodine concentration. The surface morphology of a deposited film measured by atomic force microscopy (AFM) is shown in Fig. 2. It is observed that the deposited films were having irregular surfaces with some impurity phases.

# 3.2. Structural analysis

An X-ray diffractogram of a  $(ZnSe)I_{0.003}$  film deposited on a glass substrate at room temperature is shown in Fig. 3. It is observed that the XRD pattern of the deposited film shows a more preferred orientation along the (111) plane. The (111) direction is the closepacked direction of the zinc blende structure and the deposited films are polycrystalline having the cubic zinc



Fig. 1. Rutherford backscattering spectrum of vacuum evaporated  $(ZnSe)I_{0.003}$  thin films.



Fig. 2. The surface morphology of the  $(ZnSe)I_{0.003}$  thin films deposited onto the glass substrates.

blende structure. The XRD patterns exhibit reflections at  $2\theta = 27.21^{\circ}$ , corresponding to the (111) cubic phase [7]. The particle size (D) is calculated using the Scherrer formula from the full width at half maximum ( $\beta$ ) [14]:

$$D = \frac{0.94\lambda}{\beta\cos\theta} \tag{5}$$

where  $\lambda$  is the wavelength of the X-rays used,  $\beta$  is the full width at half maximum, *D* is the particle size value and  $\theta$  is half the angle between the incident and the scattered X-rays. The dislocation density ( $\delta$ ) can be calculated using the formula [14]:

$$\delta = \frac{1}{D^2} \tag{6}$$

The strain values ( $\varepsilon$ ) can be evaluated by using the following relation [14]:

$$\varepsilon = \left[\frac{1}{D\cos\theta} - \beta\right] \times \frac{1}{\tan\theta} \tag{7}$$

From the above relations, the structural parameters such as particle size (*D*), strain ( $\varepsilon$ ) and dislocation density ( $\delta$ ) values are calculated as 28.28 nm,  $1.38 \times 10^{-3} \text{ lm}^{-2} \text{ m}^{-4}$  and  $1.29 \times 10^{15} \text{ lm/m}^2$ , respectively.

## 3.3. Spectroscopic ellipsometry

The ellipsometry angles,  $\psi$  and  $\Delta$  were directly measured by spectroscopic ellipsometry. Fitting three unknown parameters with two measured values,  $\psi$  and

 $\Delta$ , carries with it the possibility of ending up with results that lack uniqueness. However, the uniqueness of the results can be fulfilled by using the self-consistency of the Tauc–Lorentz oscillators (via the Kramers–Kroning relationship, Eq. (8)) [15]:

$$\varepsilon_1(E) = 1 + \frac{2}{\pi} \int_0^\infty \frac{E'' \varepsilon_1(E'')}{E''^2 - E^2} dE''$$
(8)

$$\varepsilon_2(E) = \frac{2E}{\pi} \int_0^\infty \frac{\varepsilon_1(E'')}{E''^2 - E^2} dE''$$
(9)

From the above relations, the value of the pseudodielectric function  $\varepsilon(E) = \varepsilon_1(E) + i\varepsilon_2(E)$  is calculated. The values of the real refractive index n(E) and extinction coefficient k(E) are calculated from the following relations [16]:

$$n(E) = \left\{ \frac{\left[\varepsilon_{1}(E)^{2} + \varepsilon_{2}(E)^{2}\right]^{1/2} + \varepsilon_{1}(E)}{2} \right\}^{1/2}$$
(10)

$$k(E) = \left\{ \frac{\left[\varepsilon_{1}(E)^{2} + \varepsilon_{2}(E)^{2}\right]^{1/2} - \varepsilon_{1}(E)}{2} \right\}^{1/2}$$
(11)

The absorption edge of ZnSe corresponds to direct transitions from the highest valence band to the lowest conduction band at the  $\Gamma$  point (i.e.,  $\Gamma_{15}^{V} \rightarrow \Gamma_{1}^{C}$ , in single-group notation). The spin-orbit interaction splits the  $\Gamma_{15}^{V}$  valence band into  $\Gamma_{8}^{V}$  and  $\Gamma_{7}^{V}$  (splitting energy,



Fig. 3. X-ray diffraction pattern of  $(ZnSe)I_{0.003}$  thin films deposited onto glass substrates.



Fig. 4. Spectral dependence of  $\varepsilon_1$  and  $\varepsilon_2$  of (ZnSe)I<sub>0.003</sub> thin films.

 $\Delta_0$  and the  $\Gamma_{15}^{C}$  conduction band into  $\Gamma_8^{C}$  and  $\Gamma_7^{C}$ (splitting energy,  $\Delta_0^{I}$ ). The corresponding optical transitions are, respectively labeled  $E_0[\Gamma_8^{V}(\Gamma_{15}^{V}) \rightarrow \Gamma_6^{C}(\Gamma_1^{C})]$ and  $E_0 + \Delta_0[\Gamma_7^{V}(\Gamma_{15}^{V}) \rightarrow \Gamma_6^{C}(\Gamma_1^{C})]$ . The spin-orbit interaction also splits the  $L_3^{V}(\Lambda_3^{V})$  valence into  $L_{4,5}^{V}(\Lambda_{4,5}^{C})$  and the  $L_3^{C}(\Lambda_3^{C})$  conduction band into  $L_6^{V}(\Lambda_6^{V})$  and  $L_{4,5}^{C}(\Lambda_{4,5}^{C})$ , respectively [17]. The corresponding optical transitions are, respectively labeled as  $E_1[L_{4,5}^{V}(L_3^{V}) \rightarrow L_6^{C}(L_1^{C})$  or  $\Lambda_{4,5}^{C}(\Lambda_3^{V}) \rightarrow \Lambda_6^{C}(\Lambda_1^{C})]$ ,  $E_1 + \Delta_1[L_6^{V}(L_3^{V}) \rightarrow L_6^{C}(L_1^{V})$  or  $\Lambda_6^{V}(\Lambda_3^{V}) \rightarrow \Lambda_6^{C}(\Lambda_1^{C})]$  [18]. These points are called critical points (CPs).

The dielectric behaviour of a crystalline material is known to be strongly connected with the energy-band structure. These CP's are associated with transitions in the energy band structure at the energies labeled as  $E_0$ ,  $E_0 + \Delta_0$ ,  $E_1$  and  $E_1 + \Delta_1$ . The  $E_0$  and  $E_0 + \Delta_0$  transitions are three dimensional (3D) M<sub>o</sub>CPs and occur in ZnSe at photon energies ~2.7 eV ( $E_0$ ) and ~3.1 eV ( $E_0 + \Delta_0$ ) at 300 K. The  $E_1$  and  $E_1 + \Delta_1$  CPs may be of the 3D M<sub>1</sub> type and occur in ZnSe at energies around ~5 eV. In



Fig. 5. Numerically calculated spectral dependence of refractive index (n) and extinction coefficient of  $(ZnSe)I_{0.003}$  thin films.



Fig. 6. Transmittance spectrum of  $(ZnSe)I_{0.003}$  thin films of 225 nm thickness deposited onto glass substrate.

Fig. 4, we can see clear peaks at 2.93 and 4.72 eV which correspond to 3D M<sub>0</sub> and 3D M<sub>1</sub> type, respectively. These transitions originate from transitions at the  $E_0$  and  $E_1$  edges. Similar results have been reported by earlier investigators [19–21]. Fig. 5 shows the spectral dependence of numerically calculated n(E) and k(E) for (ZnSe)I<sub>0.003</sub> thin films deposited onto glass substrates.

Similar to  $\varepsilon_1$  and  $\varepsilon_2$ , these spectra reveal the distinct  $E_0$  and  $E_0 + \Delta_0$  critical-point structures.

# 3.4. Optical properties

Fig. 6 shows a transmittance spectrum of a (ZnSe)  $I_{0.003}$  film deposited at room temperature. This spectrum of the film exhibits a sharp fall of transmittance at the band edge, which confirms that the deposited films have a crystalline nature. This sharp fall of transmittance occurs in the low wavelength region corresponding to



Fig. 7. Current-voltage characteristics of a Al/(ZnSe)I\_{0.003}/Al sandwich thin film capacitor.



Fig. 8. Plot between the temperature versus  $p_o$  and  $p_t$ .

the band gap of the film. The inset of Fig. 6 shows a plot between  $(\alpha h\nu)^2$  and  $h\nu$ , which yields the energy gap values of  $(ZnSe)I_{0.003}$  films. The energy band gap value of the deposited film is calculated as 2.63 eV. Similar results have been reported by Venkata et al. [22].

#### 3.5. DC conduction studies

The variation of current versus voltage graph of the prepared Al–(ZnSe)I<sub>0.003</sub>–Al device is shown in Fig. 7. In the current versus voltage characteristics, three different types of conduction mechanism were observed. At low voltages (< 2.6 V), an Ohmic conduction mechanism is observed. The voltage range between 2.65 to 5.45 V, a trap square law dependence is observed and is very clearly shown in Fig. 7. At higher voltages (> 5.5 V), i.e. in the third region for the voltage between 5.5 and 10 V, the current increases rapidly with increasing applied voltage and the slope varies from 3 to 6.43 and is seen to decrease with increasing



Fig. 9. The graph between  $\theta$  and inverse absolute temperature.

Table 1 Electrical parameters of  $(ZnSe)I_{0.003}$  thin films

| <i>T</i> (K) | $\mu \times 10^{-13} \text{ (m}^2/\text{v s)}$ | $p_0 \times 10^{22} \text{ (m}^{-3}\text{)}$ | $p_t \times 10^{22} \text{ (m}^{-3}\text{)}$ |
|--------------|--|--|--|
| 302          | 0.389  | 1.057  | 1.925  |
| 322          | 0.891  | 1.060  | 1.543  |
| 342          | 1.123  | 1.108  | 1.332  |
| 382          | 1.306  | 1.143  | 1.026  |

temperature. In the third region, the conduction mechanism is space charge limited (SCLC). The mobility value is calculated from the relation [23]:

$$J = \frac{9}{8}\mu_p \varepsilon' \varepsilon_0 \theta \frac{V^2}{d^3} \tag{12}$$

where  $\mu_p$  is the hole mobility,  $\varepsilon_0$  is the permittivity of free space and  $\varepsilon'$  is the dielectric constant of the material of the film. The value of trap density  $(p_t)$  can be calculated from the relation:

$$\frac{I_1}{I_2} = \theta = \frac{p_0}{p_0 + p_1} \tag{13}$$

The value of free carrier density  $(p_0)$  is calculated from the following relation:

$$p_0 = \frac{\varepsilon^1 \varepsilon_0 \theta}{q d^2} V_{tr} \tag{14}$$

The plot between temperature (*T*) and  $p_0$  and  $p_t$  is shown in Fig. 8. The free carrier density ( $p_0$ ) increases, but the value of trapped carrier density decreases with an increase in temperature. The relation between  $\theta$  and inverse absolute temperature is shown in Fig. 9. From the slope of the straight line graph, the activation energy has been calculated as 0.251 eV. The calculated values of carrier mobility, free carrier concentration and trap carrier density ( $p_t$ ) are given in Table 1.

### 4. Conclusions

The (ZnSe)I<sub>0.003</sub> thin films were prepared using a vacuum evaporation method. The composition of the deposited film is found to be nearly stoichiometric. In the XRD analysis, the structure of the deposited film is found to be cubic. We have measured the dielectric functions ( $\varepsilon_1$  and  $\varepsilon_2$ ) of (ZnSe)I<sub>0.003</sub> thin films grown on glass substrates using spectroscopic ellipsometry. The spectral data show three distinct critical-point structures, including the weak structures at  $E_0 + \Delta_0$ , indicating that the sample has a high crystalline quality. The optical studies revealed the existence of direct allowed transitions in the deposited films. A Space charge limited conduction (SCLC) mechanism is observed from the d.c. conduction studies.

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